

**substance: titanium oxide (TiO<sub>2</sub>)**

**property: properties of vibrational and Raman modes in rutile**

**symmetry of optical modes: Fig. 1**

**phonon dispersion curves: Figs. 2, 3.**

**frequencies of optical lattice modes**

(data in 10<sup>12</sup> Hz)

$\nu(\Gamma_4^+, B_{2g})$	24.78	R, $T = 4$ K	R: Raman data from [67P]
	24.72(25)	N	I: IR data from [64E]
$\nu(\Gamma_5^+, E_u)_{LO}$	24.18	I	N: neutron data from [71T]
	25.24 (34)	N	Symmetries of modes are given in
$\nu(\Gamma_1^-, A_{2u})_{LO}$	24.33	I	Koster notation (as in Fig. 2)
	24.30	N	and Mulliken notation.
$\nu(\Gamma_1^+, A_{1g})$	18.36	R	For the correspondence of the
	18.30	N	data between $\Gamma$ , X, M and Z
$\nu(\Gamma_5^+, E_u)_{TO}$	15.00	I	points, see Figs. 2, 3
	14.81(15)	N	
$\nu(\Gamma_5^-, E_g)$	13.41	R	
	13.339(181)	N	
$\nu(\Gamma_5^+, E_u)_{LO}$	13.74	I	
	12.853(129)	N	
$\nu(\Gamma_4^-, B_{1u}^2)$	12.182(122)	N	
$\nu(\Gamma_5^+, E_u)_{TO}$	11.64	I	
$\nu(\Gamma_5^+, E_u)_{LO}$	11.19	I	
	11.232(112)	N	
$\nu(\Gamma_5^+, E_u)_{TO}$	5.49	I	
	5.661(75)	N	
$\nu(\Gamma_1^-, A_{2u})_{TO}$	5.01	I	
	5.177(52)	N	
$\nu(\Gamma_3^+, B_{1g})$	4.29	R	
	4.246(94)	N	
$\nu(\Gamma_4^-, B_{1u})$	3.389(57)	N	
$\nu(X_1)$	25.38(25)		corresponding to
$\nu(X_1)$	24.38(25)		( $\Gamma_4^+$ )
$\nu(X_2)$	13.528(135)		( $\Gamma_5^+$ )
$\nu(X_1)$	11.756(118)		( $\Gamma_5^-$ )
$\nu(X_1)$	9.130(125)		( $\Gamma_5^+$ )
$\nu(X_2)$	9.033(90)		( $\Gamma_5^+$ )
$\nu(X_1)$	8.071(81)		( $\Gamma_4^-$ ) and ( $\Gamma_1^-$ )
$\nu(X_1)$	5.821(58)		( $\Gamma_5^+$ ) and ( $\Gamma_3^+$ )
$\nu(X_2)$	3.084(31)		acoustic LA and TA branches ( $\Gamma_5^+$ )
$\nu(M_{1,2})$	23.26(23)		( $\Gamma_4^-$ ) and acoustic TA branch ( $\Gamma_1^-$ )
$\nu(M_9^+)$	13.522(135)		( $\Gamma_4^+$ ) and ( $\Gamma_5^+$ )
$\nu(M_9^+)$	9.442(94)		( $\Gamma_5^-$ )
$\nu(M_{5,6})$	9.398(94)		( $\Gamma_5^+$ )
$\nu(M_9^+)$	7.811(78)		( $\Gamma_4^-$ ) and ( $\Gamma_1^-$ )
$\nu(M_{5,6})$	3.058(56)		( $\Gamma_5^+$ ) and ( $\Gamma_3^+$ )
$\nu(M_9^+)$	2.936(70)		( $\Gamma_4^-$ ) and acoustic TA branch ( $\Gamma_1^-$ )
$\nu(Z_1)$	12.397(124)		acoustic LA and TA branches ( $\Gamma_5^+$ )
$\nu(Z_2)$	12.312(123)		( $\Gamma_4^-$ ) and acoustic LA-branch ( $\Gamma_1^-$ )
			( $\Gamma_5^-$ )

71T

$\nu(Z_4)$	9.720(97)	$(\Gamma_5^+)$
$\nu(Z_3)$	9.040(90)	$(\Gamma_3^+)$
$\nu(Z_2)$	3.697 (69)	acoustic TA branch $(\Gamma_5^+)$

**temperature dependence of vibrational modes:**  $A_{2u}$  and  $E_u(TO)$  modes, with displacements along  $[001]$  and  $[110]$ , have an anomalous increase in intensity and decrease in frequency with decreasing temperature (Fig. 4) characteristic for mode softening on the approach to a ferroelectric catastrophe. These modes are responsible for the high dielectric constant.

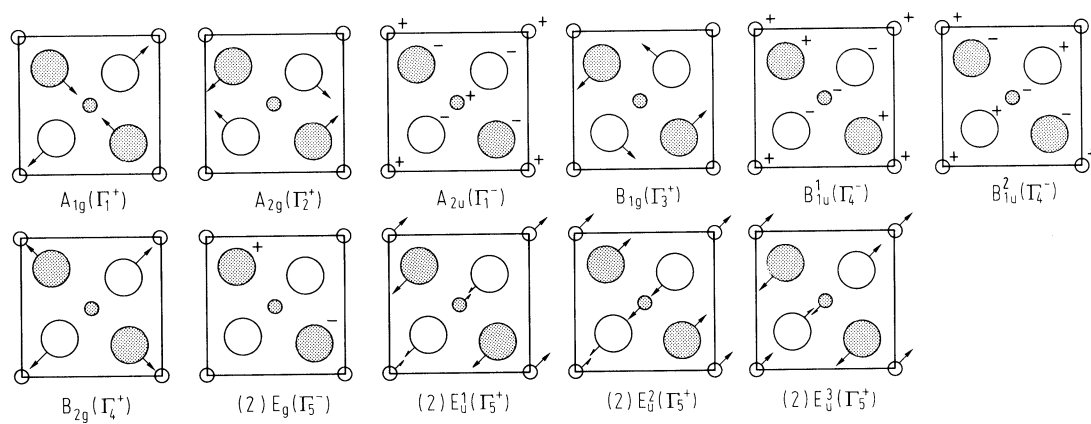
**pressure dependence of Raman modes:**  $E_g$  ( $417\text{ cm}^{-1}$ ) and  $A_{1g}$  ( $612\text{ cm}^{-1}$ ) peaks shift with rates of  $+0.43$  and  $+0.41\text{ cm}^{-1}\text{ kbar}^{-1}$ , but the  $B_{1g}$  ( $145\text{ cm}^{-1}$ ) mode softens by  $-0.36\text{ cm}^{-1}\text{ kbar}^{-1}$  [71N]. Above 26 kbar a transition to a new spectrum is observed [71N, 79H], which has been interpreted as a transition from an anharmonic to a harmonic material at high pressure [79H] and as a transition to a new phase of  $\text{PbO}_2$  structure [71N, 80M] or  $\text{CaCl}_2$  structure [78N]. Uniaxial stress measurements [80M] show that all modes displace linearly with pressure, but the  $B_{1g}$  mode shows non-classical behaviour. At very high pressures ( $> 300\text{ kbar}$ ) another phase is observed in Raman spectrum whose structure is unknown [81M].

## References:

- 61P     Parker, R. A.: Phys. Rev. 124 (1961) 1719.
- 64E     Eagles, D. M.: J. Phys. Chem. Solids 25 (1964) 1243.
- 67P     Porto, S. P. S., Fleury, P. A., Damen, T. C.: Phys. Rev. 154 (1967) 522.
- 71N     Nicol, M., Fong, M. Y.: J. Chem. Phys. 54 (1971) 3167.
- 71T     Traylor, J. G., Smith, H. G., Nicklow, R. M., Wilkinson, M. K.: Phys. Rev. B3 (1971) 3457.
- 78N     Nicol, M., Hara, Y.: Proc. Int. Conf. Raman Spectrosc. 6th 1978, 2 320; E. D. Schmid, R. S. Krishnan and W. Kiefer (eds.) London: Heyden.
- 79H     Hara, Y., Nicol, M.: Phys. Status Solidi (b) 94 (1979) 317.
- 80M     Merle, P., Pascual, J., Camassel, J., Mathieu, H.: Phys. Rev. B21 (1980) 1617.
- 81M     Mammone, J. F., Nicol, M., Sharma, S. K.: J. Phys. Chem. Solids 42 (1981) 379.

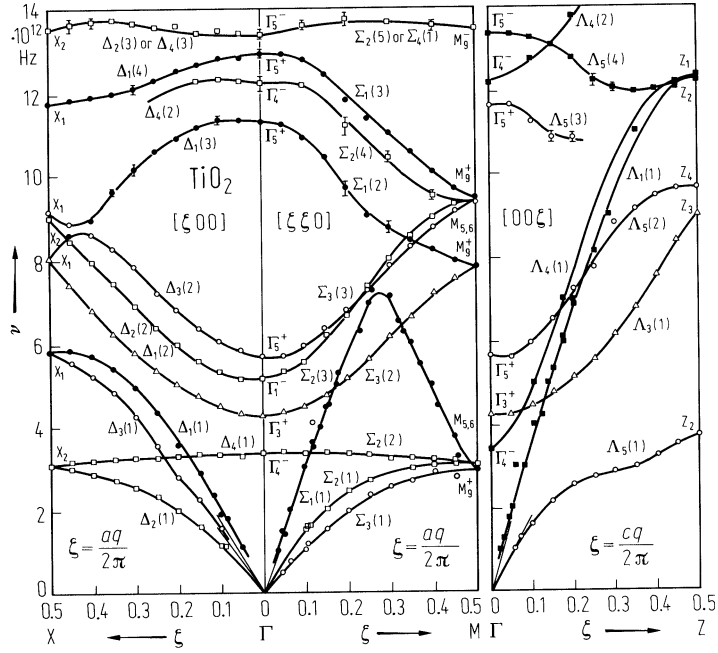
**Fig. 1.**

TiO<sub>2</sub>. Symmetry of the optical modes in rutile at  $q = 0$ . Projection drawn down the  $c$ -axis [71T]



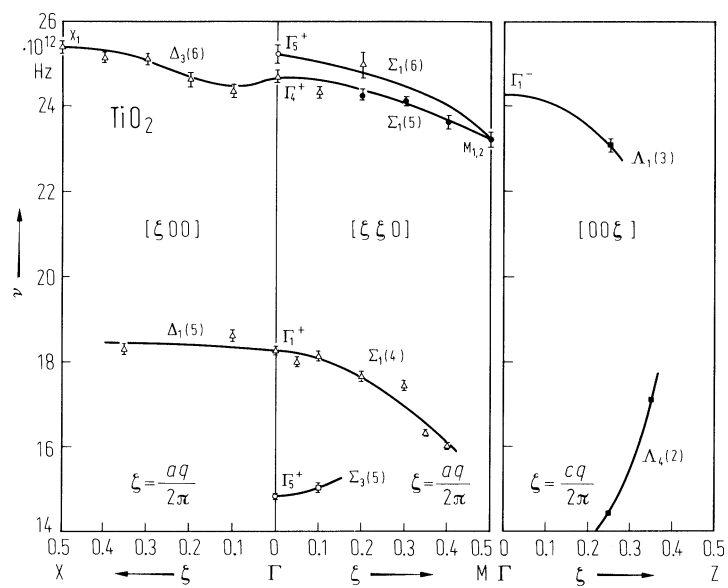
**Fig. 2.**

$\text{TiO}_2$ . Low-frequency phonon dispersion curves for rutile. Lines through the points are guide lines only. Solid (open) points represent measurements made under predominantly longitudinal (transverse) scattering conditions. Circular (square) points indicate that the scattering vector  $q$  lay predominantly normal to (parallel to) the  $z$ -axis. Triangular points are modes with displacements normal to  $z$ , but which may not be classified as longitudinal or transverse [71T].



**Fig. 3.**

TiO<sub>2</sub>. High-frequency phonon dispersion spectrum. Notation as in Fig. 2 [71T].



**Fig. 4.**

TiO<sub>2</sub>. Square of the A<sub>2u</sub>-mode frequency vs. temperature. Solid line calculated from dielectric measurements [61P], circles from neutron scattering [71T].

